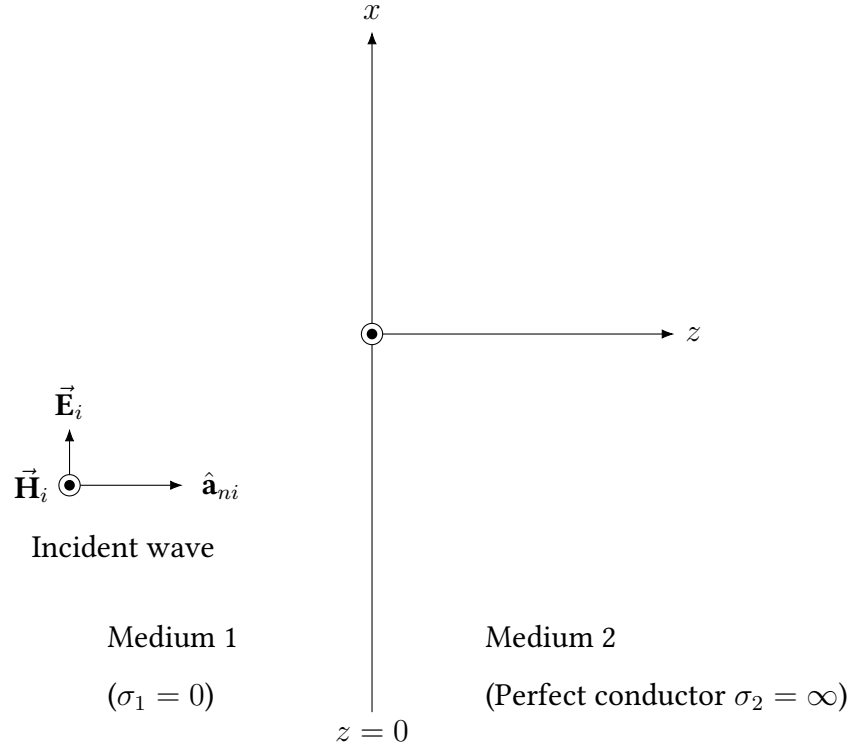


3. Normal Incidence at a Plane Conducting Boundary



$$\vec{\mathbf{E}}_i(z) = \hat{\mathbf{a}}_x E_{i0} e^{-j\beta_1 z} \quad (1)$$

$$\vec{\mathbf{H}}_i(z) = \hat{\mathbf{a}}_y \frac{E_{i0}}{\eta_1} e^{-j\beta_1 z} \quad (2)$$

E_{i0} : the magnitude of $\vec{\mathbf{E}}_i$ at $z = 0$

β_1 : phase constant of medium 1

η_1 : intrinsic impedance of medium 1

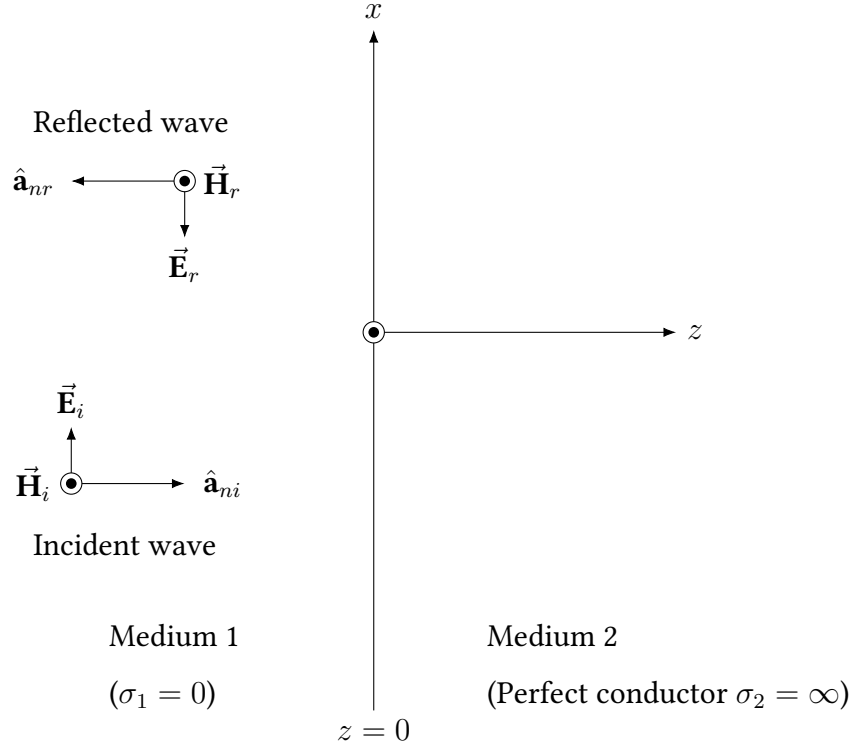
$$\eta_1 = \sqrt{\frac{\mu_1}{\epsilon_1}} \quad (3)$$

In a perfect conductor electric and magnetic fields are zero.

$$\vec{\mathbf{E}}_2 = 0 \quad (4)$$

$$\vec{\mathbf{H}}_2 = 0 \quad (5)$$

The incident wave is reflected giving rise to a reflected wave ($\vec{\mathbf{E}}_r, \vec{\mathbf{H}}_r$).



$$\vec{\mathbf{E}}_r(z) = -\hat{\mathbf{a}}_x E_{r0} e^{j\beta_1 z} \quad (6)$$

Reflected wave travels in $-z$ direction.

Total field in medium 1

$$\vec{\mathbf{E}}_1(z) = \vec{\mathbf{E}}_i(z) + \vec{\mathbf{E}}_r(z) \quad (7)$$

$$\vec{\mathbf{E}}_1(z) = \hat{\mathbf{a}}_x E_{i0} e^{-j\beta_1 z} - \hat{\mathbf{a}}_x E_{r0} e^{j\beta_1 z} \quad (8)$$

$$\vec{\mathbf{E}}_1(z) = \hat{\mathbf{a}}_x (E_{i0} e^{-j\beta_1 z} - E_{r0} e^{j\beta_1 z}) \quad (9)$$

$$\vec{\mathbf{E}}_2 = 0 \quad (10)$$

$$E_{1t} = E_{2t} = 0 \quad (11)$$

$$\vec{\mathbf{E}}_1(0) = 0 = \hat{\mathbf{a}}_x (E_{i0} - E_{r0}) \quad (12)$$

$$\Rightarrow \boxed{E_{r0} = E_{i0}} \quad (13)$$

$$\vec{\mathbf{E}}_r(z) = -\hat{\mathbf{a}}_x E_{i0} e^{j\beta_1 z} \quad (14)$$

$$\vec{\mathbf{E}}_1(z) = \hat{\mathbf{a}}_x E_{i0} (e^{-j\beta_1 z} - e^{j\beta_1 z}) \quad (15)$$

$$\vec{\mathbf{E}}_1(z) = \hat{\mathbf{a}}_x E_{i0} (-2j) \sin \beta_1 z \quad (16)$$

$$\vec{\mathbf{E}}_1(z) = -\hat{\mathbf{a}}_x 2j E_{i0} \sin \beta_1 z \quad (17)$$

$$\vec{\mathbf{H}}_r(z) = \frac{1}{\eta_1} \hat{\mathbf{a}}_{nr} \times \vec{\mathbf{E}}_r(z) \quad (18)$$

$$\hat{\mathbf{a}}_{nr} = -\hat{\mathbf{a}}_z \quad (19)$$

$$\vec{\mathbf{H}}_r(z) = \frac{1}{\eta_1} (-\hat{\mathbf{a}}_z) \times (-\hat{\mathbf{a}}_x) E_{i0} e^{j\beta_1 z} \quad (20)$$

$$\vec{\mathbf{H}}_r(z) = \frac{1}{\eta_1} \hat{\mathbf{a}}_y E_{i0} e^{j\beta_1 z} \quad (21)$$

$$\vec{\mathbf{H}}_1(z) = \vec{\mathbf{H}}_i(z) + \vec{\mathbf{H}}_r(z) \quad (22)$$

$$\vec{\mathbf{H}}_1(z) = \hat{\mathbf{a}}_y \frac{E_{i0}}{\eta_1} e^{-j\beta_1 z} + \hat{\mathbf{a}}_y \frac{E_{i0}}{\eta_1} e^{j\beta_1 z} \quad (23)$$

$$\vec{\mathbf{H}}_1(z) = \hat{\mathbf{a}}_y \frac{E_{i0}}{\eta_1} (e^{-j\beta_1 z} + e^{j\beta_1 z}) \quad (24)$$

$$\vec{\mathbf{H}}_1(z) = \hat{\mathbf{a}}_y 2 \frac{E_{i0}}{\eta_1} \cos \beta_1 z \quad (25)$$

$$\vec{\mathbf{E}}_1(z, t) = \text{Re} \left[\vec{\mathbf{E}}_1(z) e^{j\omega t} \right] \quad (26)$$

$$\vec{\mathbf{E}}_1(z, t) = \text{Re} \left[-\hat{\mathbf{a}}_x 2j E_{i0} \sin \beta_1 z e^{j\omega t} \right] \quad (27)$$

$$\vec{\mathbf{E}}_1(z, t) = \text{Re} \left[-\hat{\mathbf{a}}_x 2j E_{i0} \sin \beta_1 z (\cos \omega t + j \sin \omega t) \right] \quad (28)$$

$$\vec{\mathbf{E}}_1(z, t) = \hat{\mathbf{a}}_x 2 E_{i0} \sin \beta_1 z \sin \omega t \quad (29)$$

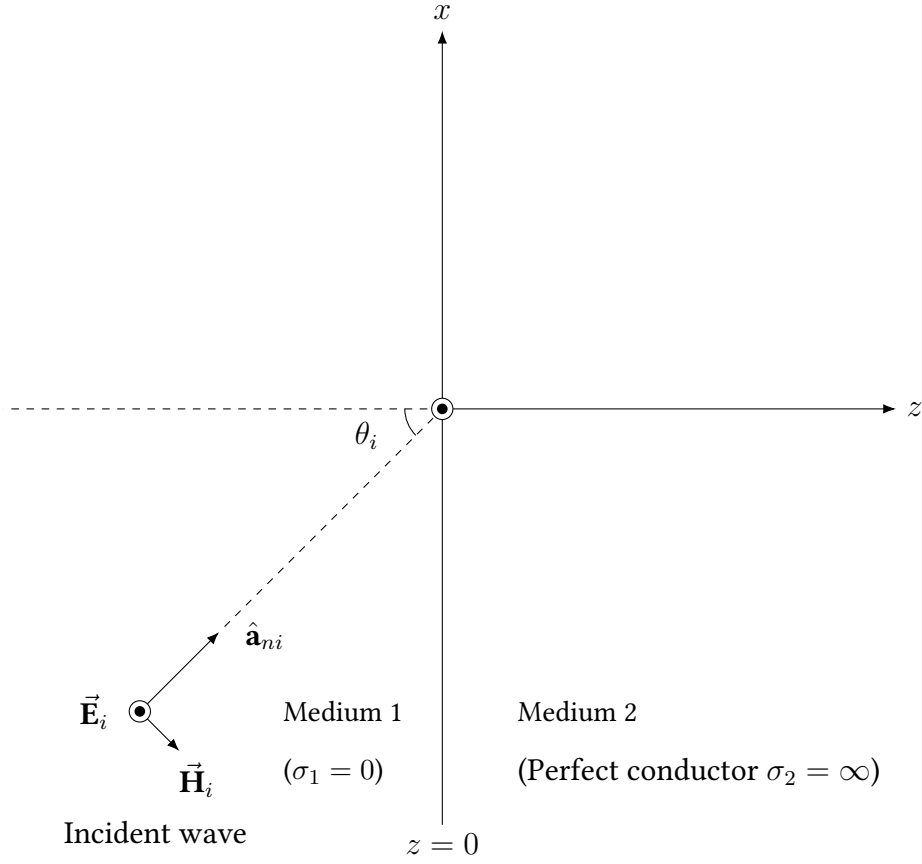
$$\vec{\mathbf{H}}_1(z, t) = \text{Re} \left[\vec{\mathbf{H}}_1(z) e^{j\omega t} \right] \quad (30)$$

$$\vec{\mathbf{H}}_1(z, t) = \text{Re} \left[\hat{\mathbf{a}}_y 2 \frac{E_{i0}}{\eta_1} \cos \beta_1 z e^{j\omega t} \right] \quad (31)$$

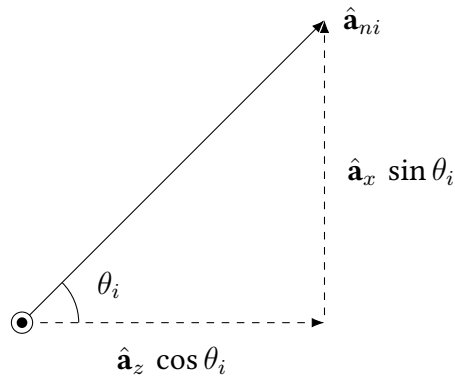
$$\vec{\mathbf{H}}_1(z, t) = \hat{\mathbf{a}}_y 2 \frac{E_{i0}}{\eta_1} \cos \beta_1 z \cos \omega t \quad (32)$$

4. Oblique Incidence at a Plane Conducting Boundary

4.1 Horizontal Polarization (Yatay Kutuplanma)



θ_i : angle of incidence



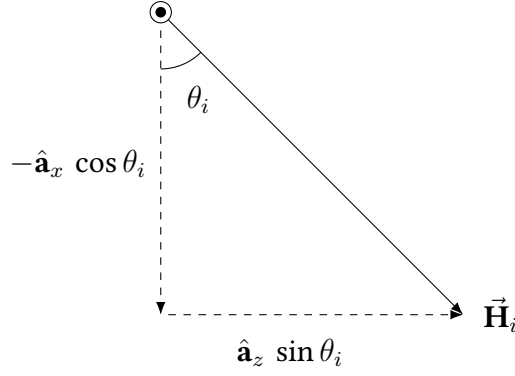
$$\hat{\mathbf{a}}_{ni} = \hat{\mathbf{a}}_x \sin \theta_i + \hat{\mathbf{a}}_z \cos \theta_i \quad (33)$$

$$\vec{\mathbf{E}}_i(x, z) = \hat{\mathbf{a}}_y E_{i0} e^{-j\beta_1 \hat{\mathbf{a}}_{ni} \cdot \vec{\mathbf{R}}} \quad (34)$$

$$\hat{\mathbf{a}}_{ni} \cdot \vec{\mathbf{R}} = (\hat{\mathbf{a}}_x \sin \theta_i + \hat{\mathbf{a}}_z \cos \theta_i) \cdot (\hat{\mathbf{a}}_x x + \hat{\mathbf{a}}_y y + \hat{\mathbf{a}}_z z) \quad (35)$$

$$\hat{\mathbf{a}}_{ni} \cdot \vec{\mathbf{R}} = (x \sin \theta_i + z \cos \theta_i) \quad (36)$$

$$\vec{\mathbf{E}}_i(x, z) = \hat{\mathbf{a}}_y E_{i0} e^{-j\beta_1 (x \sin \theta_i + z \cos \theta_i)} \quad (37)$$

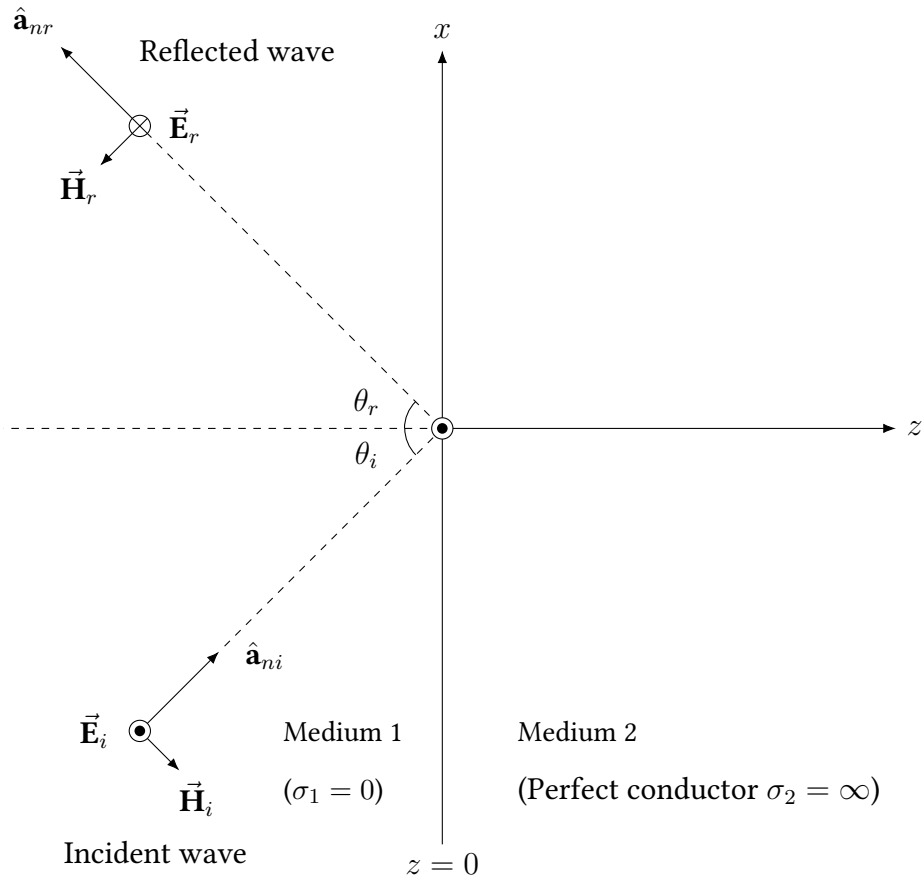


$$\vec{\mathbf{H}}_i(x, z) = \frac{1}{\eta_1} \left[\hat{\mathbf{a}}_{ni} \times \vec{\mathbf{E}}_i(x, z) \right] \quad (38)$$

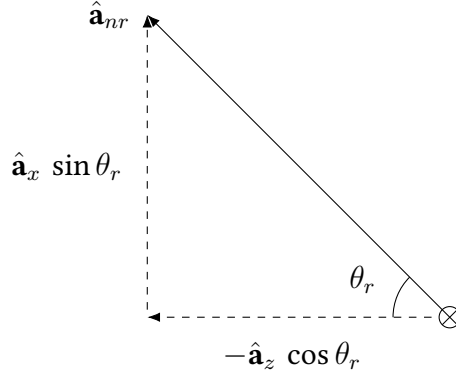
$$\vec{\mathbf{H}}_i(x, z) = \frac{1}{\eta_1} \left[(\hat{\mathbf{a}}_x \sin \theta_i + \hat{\mathbf{a}}_z \cos \theta_i) \times \hat{\mathbf{a}}_y E_{i0} e^{-j\beta_1(x \sin \theta_i + z \cos \theta_i)} \right] \quad (39)$$

$$\vec{\mathbf{H}}_i(x, z) = (\hat{\mathbf{a}}_z \sin \theta_i - \hat{\mathbf{a}}_x \cos \theta_i) \frac{E_{i0}}{\eta_1} e^{-j\beta_1(x \sin \theta_i + z \cos \theta_i)} \quad (40)$$

$$\vec{\mathbf{H}}_i(x, z) = (-\hat{\mathbf{a}}_x \cos \theta_i + \hat{\mathbf{a}}_z \sin \theta_i) \frac{E_{i0}}{\eta_1} e^{-j\beta_1(x \sin \theta_i + z \cos \theta_i)} \quad (41)$$



θ_r : angle of reflection



$$\vec{\mathbf{E}}_r(x, z) = -\hat{\mathbf{a}}_y E_{r0} e^{-j\beta_1 \hat{\mathbf{a}}_{nr} \cdot \vec{\mathbf{R}}} \quad (42)$$

$$\hat{\mathbf{a}}_{nr} = \hat{\mathbf{a}}_x \sin \theta_r - \hat{\mathbf{a}}_z \cos \theta_r \quad (43)$$

$$\hat{\mathbf{a}}_{nr} \cdot \vec{\mathbf{R}} = (\hat{\mathbf{a}}_x \sin \theta_r - \hat{\mathbf{a}}_z \cos \theta_r) \cdot (\hat{\mathbf{a}}_x x + \hat{\mathbf{a}}_y y + \hat{\mathbf{a}}_z z) \quad (44)$$

$$\hat{\mathbf{a}}_{nr} \cdot \vec{\mathbf{R}} = (x \sin \theta_r - z \cos \theta_r) \quad (45)$$

$$\vec{\mathbf{E}}_r(x, z) = -\hat{\mathbf{a}}_y E_{r0} e^{-j\beta_1 (x \sin \theta_r - z \cos \theta_r)} \quad (46)$$

At the boundary surface, $z = 0$, the tangential components of $\vec{\mathbf{E}}_1$ is zero ($E_{1t} = E_{2t} = 0$).

$$\vec{\mathbf{E}}_1(x, z) = \vec{\mathbf{E}}_i(x, z) + \vec{\mathbf{E}}_r(x, z) \quad (47)$$

$$\vec{\mathbf{E}}_1(x, z = 0) = \vec{\mathbf{E}}_i(x, z = 0) + \vec{\mathbf{E}}_r(x, z = 0) = 0 \quad (48)$$

$$\vec{\mathbf{E}}_1(x, z = 0) = \hat{\mathbf{a}}_y E_{i0} e^{-j\beta_1 x \sin \theta_i} - \hat{\mathbf{a}}_y E_{r0} e^{-j\beta_1 x \sin \theta_r} = 0 \quad (49)$$

$$E_{i0} e^{-j\beta_1 x \sin \theta_i} = E_{r0} e^{-j\beta_1 x \sin \theta_r} \quad \text{valid for all } x \quad (50)$$

For $x = 0$,

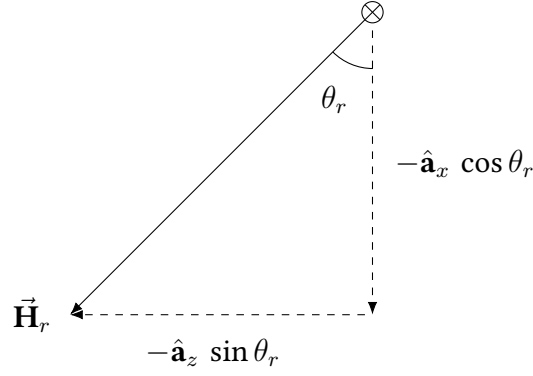
$$E_{r0} = E_{i0} \quad (51)$$

$$\Rightarrow \theta_r = \theta_i \quad (\text{Snell's law of reflection}) \quad (52)$$

The angle of reflection is equal to the angle of incidence. So we have

$$\vec{\mathbf{E}}_r(x, z) = -\hat{\mathbf{a}}_y E_{i0} e^{-j\beta_1 (x \sin \theta_i - z \cos \theta_i)} \quad (53)$$

$$\vec{\mathbf{H}}_r(x, z) = \frac{1}{\eta_1} [\hat{\mathbf{a}}_{nr} \times \vec{\mathbf{E}}_r(x, z)] \quad (54)$$



$$\vec{H}_r(x, z) = \frac{1}{\eta_1} [(\hat{a}_x \sin \theta_r - \hat{a}_z \cos \theta_r) \times (-) \hat{a}_y E_{i0} e^{-j\beta_1(x \sin \theta_i - z \cos \theta_i)}] \quad (55)$$

$$\vec{H}_r(x, z) = (-\hat{a}_z \sin \theta_r - \hat{a}_x \cos \theta_r) \frac{E_{i0}}{\eta_1} e^{-j\beta_1(x \sin \theta_i - z \cos \theta_i)} \quad (56)$$

$$\vec{H}_r(x, z) = (-\hat{a}_x \cos \theta_i - \hat{a}_z \sin \theta_i) \frac{E_{i0}}{\eta_1} e^{-j\beta_1(x \sin \theta_i - z \cos \theta_i)} \quad (57)$$

The total electric field

$$\vec{E}_1(x, z) = \vec{E}_i(x, z) + \vec{E}_r(x, z) \quad (58)$$

$$\vec{E}_1(x, z) = \hat{a}_y E_{i0} e^{-j\beta_1(x \sin \theta_i + z \cos \theta_i)} - \hat{a}_y E_{i0} e^{-j\beta_1(x \sin \theta_i - z \cos \theta_i)} \quad (59)$$

$$\vec{E}_1(x, z) = \hat{a}_y E_{i0} (e^{-j\beta_1 z \cos \theta_i} - e^{j\beta_1 z \cos \theta_i}) e^{-j\beta_1 x \sin \theta_i} \quad (60)$$

$$e^{-j\beta_1 z \cos \theta_i} - e^{j\beta_1 z \cos \theta_i} = -2j \sin(\beta_1 z \cos \theta_i) \quad (61)$$

$$\vec{E}_1(x, z) = -\hat{a}_y j 2 E_{i0} \sin(\beta_1 z \cos \theta_i) e^{-j\beta_1 x \sin \theta_i} \quad (62)$$

The total magnetic field

$$\vec{H}_1(x, z) = \vec{H}_i(x, z) + \vec{H}_r(x, z) \quad (63)$$

$$\begin{aligned} \vec{H}_1(x, z) &= (-\hat{a}_x \cos \theta_i + \hat{a}_z \sin \theta_i) \frac{E_{i0}}{\eta_1} e^{-j\beta_1(x \sin \theta_i + z \cos \theta_i)} \\ &+ (-\hat{a}_x \cos \theta_i - \hat{a}_z \sin \theta_i) \frac{E_{i0}}{\eta_1} e^{-j\beta_1(x \sin \theta_i - z \cos \theta_i)} \end{aligned} \quad (64)$$

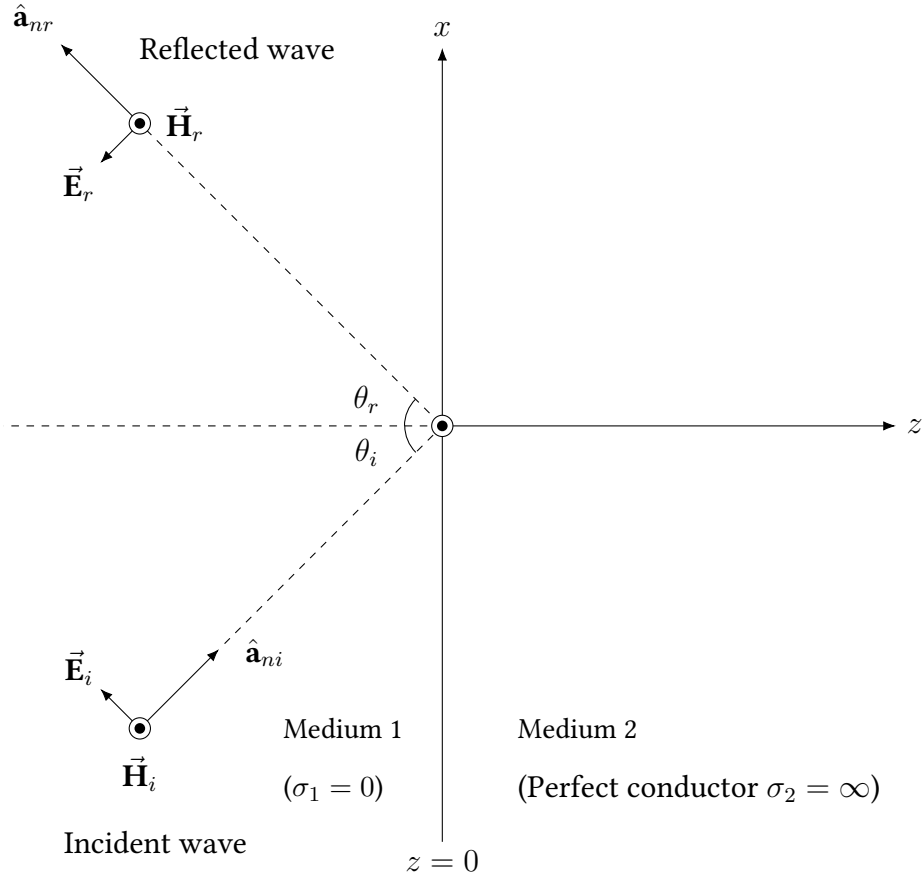
$$\begin{aligned} \vec{H}_1(x, z) &= -\hat{a}_x \cos \theta_i \frac{E_{i0}}{\eta_1} (e^{-j\beta_1 z \cos \theta_i} + e^{j\beta_1 z \cos \theta_i}) e^{-j\beta_1 x \sin \theta_i} \\ &+ \hat{a}_z \sin \theta_i \frac{E_{i0}}{\eta_1} (e^{-j\beta_1 z \cos \theta_i} - e^{j\beta_1 z \cos \theta_i}) e^{-j\beta_1 x \sin \theta_i} \end{aligned} \quad (65)$$

$$e^{-j\beta_1 z \cos \theta_i} + e^{j\beta_1 z \cos \theta_i} = 2 \cos(\beta_1 z \cos \theta_i) \quad (66)$$

$$e^{-j\beta_1 z \cos \theta_i} - e^{j\beta_1 z \cos \theta_i} = -2j \sin(\beta_1 z \cos \theta_i) \quad (67)$$

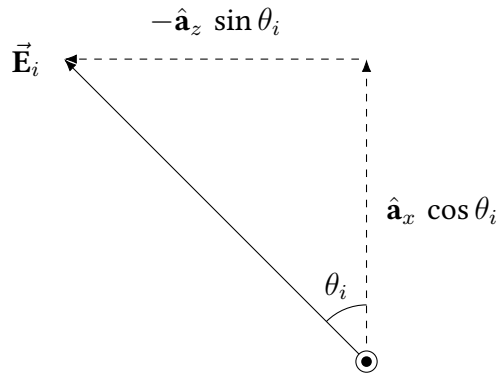
$$\begin{aligned} \vec{H}_1(x, z) &= -\hat{a}_x \frac{2E_{i0}}{\eta_1} \cos \theta_i \cos(\beta_1 z \cos \theta_i) e^{-j\beta_1 x \sin \theta_i} \\ &- \hat{a}_z \frac{2jE_{i0}}{\eta_1} \sin \theta_i \sin(\beta_1 z \cos \theta_i) e^{-j\beta_1 x \sin \theta_i} \end{aligned} \quad (68)$$

4.2 Vertical Polarization (Dikey Kutuplanma)



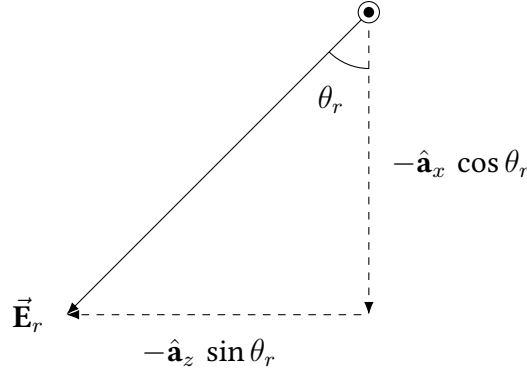
$$\hat{\mathbf{a}}_{ni} = \hat{\mathbf{a}}_x \sin \theta_i + \hat{\mathbf{a}}_z \cos \theta_i \quad (69)$$

$$\hat{\mathbf{a}}_{nr} = \hat{\mathbf{a}}_x \sin \theta_r - \hat{\mathbf{a}}_z \cos \theta_r \quad (70)$$



$$\vec{\mathbf{E}}_i(x, z) = (\hat{\mathbf{a}}_x \cos \theta_i - \hat{\mathbf{a}}_z \sin \theta_i) E_{i0} e^{-j\beta_1(x \sin \theta_i + z \cos \theta_i)} \quad (71)$$

$$\vec{\mathbf{H}}_i(x, z) = \hat{\mathbf{a}}_y \frac{E_{i0}}{\eta_1} e^{-j\beta_1(x \sin \theta_i + z \cos \theta_i)} \quad (72)$$



$$\vec{E}_r(x, z) = (-\hat{a}_x \cos \theta_r - \hat{a}_z \sin \theta_r) E_{r0} e^{-j\beta_1(x \sin \theta_r - z \cos \theta_r)} \quad (73)$$

$$\vec{H}_r(x, z) = \hat{a}_y \frac{E_{r0}}{\eta_1} e^{-j\beta_1(x \sin \theta_r - z \cos \theta_r)} \quad (74)$$

At $z = 0$, the tangential component of the total electric field intensity must vanish.

$$E_{ix}(x, z = 0) + E_{rx}(x, z = 0) = 0 \quad (75)$$

$$E_{i0} \cos \theta_i e^{-j\beta_1 x \sin \theta_i} - E_{r0} \cos \theta_r e^{-j\beta_1 x \sin \theta_r} = 0 \quad (76)$$

$$\Rightarrow E_{r0} = E_{i0} \quad (77)$$

$$\theta_r = \theta_i \quad (\text{Snell's law of reflection}) \quad (78)$$

The total electric field intensity in medium 1

$$\vec{E}_1(x, z) = \vec{E}_i(x, z) + \vec{E}_r(x, z) \quad (79)$$

$$\begin{aligned} \vec{E}_1(x, z) = & (\hat{a}_x \cos \theta_i - \hat{a}_z \sin \theta_i) E_{i0} e^{-j\beta_1(x \sin \theta_i + z \cos \theta_i)} \\ & - (\hat{a}_x \cos \theta_i + \hat{a}_z \sin \theta_i) E_{i0} e^{-j\beta_1(x \sin \theta_i - z \cos \theta_i)} \end{aligned} \quad (80)$$

$$\begin{aligned} \vec{E}_1(x, z) = & \hat{a}_x \cos \theta_i E_{i0} (e^{-j\beta_1 z \cos \theta_i} - e^{j\beta_1 z \cos \theta_i}) e^{-j\beta_1 x \sin \theta_i} \\ & - \hat{a}_z \sin \theta_i E_{i0} (e^{-j\beta_1 z \cos \theta_i} + e^{j\beta_1 z \cos \theta_i}) e^{-j\beta_1 x \sin \theta_i} \end{aligned} \quad (81)$$

$$e^{-j\beta_1 z \cos \theta_i} - e^{j\beta_1 z \cos \theta_i} = -2j \sin(\beta_1 z \cos \theta_i) \quad (82)$$

$$e^{-j\beta_1 z \cos \theta_i} + e^{j\beta_1 z \cos \theta_i} = 2 \cos(\beta_1 z \cos \theta_i) \quad (83)$$

$$\begin{aligned} \vec{E}_1(x, z) = & -\hat{a}_x 2j E_{i0} \cos \theta_i \sin(\beta_1 z \cos \theta_i) e^{-j\beta_1 x \sin \theta_i} \\ & - \hat{a}_z 2 E_{i0} \sin \theta_i \cos(\beta_1 z \cos \theta_i) e^{-j\beta_1 x \sin \theta_i} \end{aligned} \quad (84)$$

Total magnetic field in medium 1

$$\vec{H}_1(x, z) = \vec{H}_i(x, z) + \vec{H}_r(x, z) \quad (85)$$

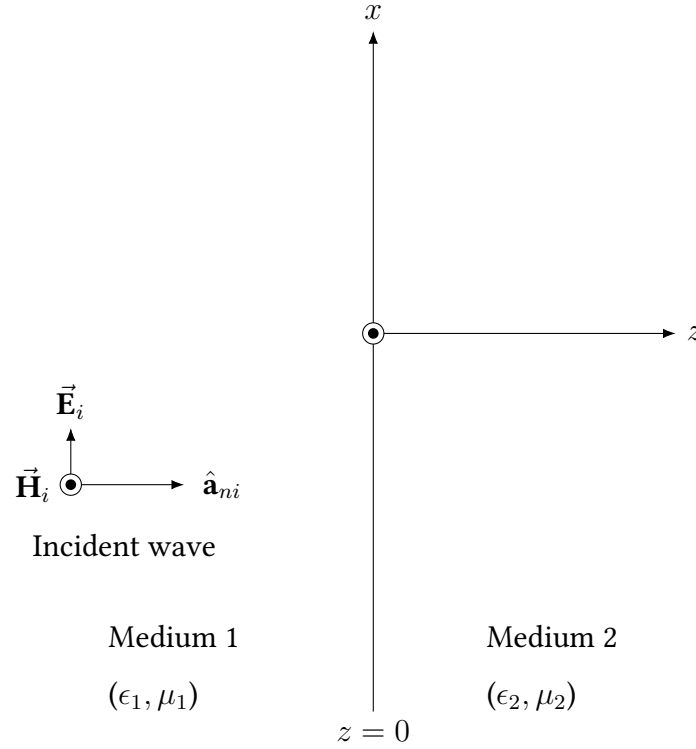
$$\vec{H}_1(x, z) = \hat{a}_y \frac{E_{i0}}{\eta_1} e^{-j\beta_1(x \sin \theta_i + z \cos \theta_i)} + \hat{a}_y \frac{E_{i0}}{\eta_1} e^{-j\beta_1(x \sin \theta_i - z \cos \theta_i)} \quad (86)$$

$$\vec{\mathbf{H}}_1(x, z) = \hat{\mathbf{a}}_y \frac{E_{i0}}{\eta_1} \left(e^{-j\beta_1 z \cos \theta_i} + e^{j\beta_1 z \cos \theta_i} \right) e^{-j\beta_1 x \sin \theta_i} \quad (87)$$

$$e^{-j\beta_1 z \cos \theta_i} + e^{j\beta_1 z \cos \theta_i} = 2 \cos(\beta_1 z \cos \theta_i) \quad (88)$$

$$\vec{\mathbf{H}}_1(x, z) = \hat{\mathbf{a}}_y \frac{2E_{i0}}{\eta_1} \cos(\beta_1 z \cos \theta_i) e^{-j\beta_1 x \sin \theta_i} \quad (89)$$

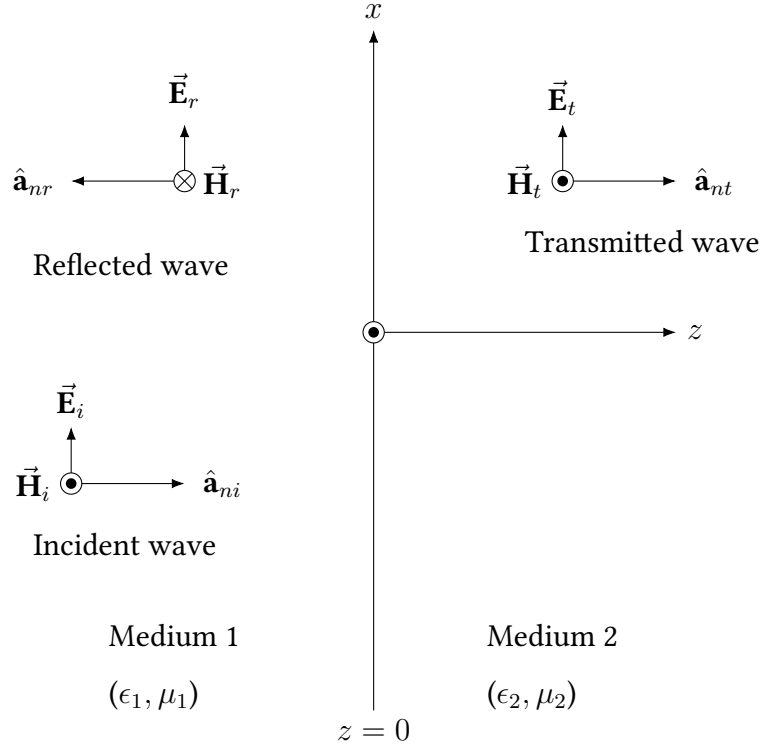
5. Normal Incidence at a Plane Dielectric Boundary



$$\vec{\mathbf{E}}_i(z) = \hat{\mathbf{a}}_x E_{i0} e^{-j\beta_1 z} \quad (90)$$

$$\vec{\mathbf{H}}_i(z) = \hat{\mathbf{a}}_y \frac{E_{i0}}{\eta_1} e^{-j\beta_1 z} \quad (91)$$

When an electromagnetic wave is incident on the surface of a different dielectric medium, part of the incident power is reflected and part is transmitted.



$$\vec{\mathbf{E}}_r(z) = \hat{\mathbf{a}}_x E_{r0} e^{j\beta_1 z} \quad (92)$$

$$\vec{\mathbf{H}}_r(z) = \frac{1}{\eta_1} (-\hat{\mathbf{a}}_z) \times \vec{\mathbf{E}}_r(z) = -\hat{\mathbf{a}}_y \frac{E_{r0}}{\eta_1} e^{j\beta_1 z} \quad (93)$$

$$\vec{\mathbf{E}}_t(z) = \hat{\mathbf{a}}_x E_{t0} e^{-j\beta_2 z} \quad (94)$$

$$\vec{\mathbf{H}}_t(z) = \frac{1}{\eta_2} \hat{\mathbf{a}}_z \times \vec{\mathbf{E}}_t(z) = \hat{\mathbf{a}}_y \frac{E_{t0}}{\eta_2} e^{-j\beta_2 z} \quad (95)$$

E_{t0} : the magnitude of $\vec{\mathbf{E}}_t$ at $z = 0$.

β_2 : phase constant of medium 2.

η_2 : intrinsic impedance of medium 2.

$$\beta_2 = \omega \sqrt{\mu_2 \epsilon_2} \quad (96)$$

$$\eta_2 = \sqrt{\frac{\mu_2}{\epsilon_2}} \quad (97)$$

E_{r0} and E_{t0} may be positive or negative, depending on $\epsilon_1, \mu_1, \epsilon_2, \mu_2$. E_{r0} and E_{t0} must be determined.

The tangential components (the x - and y - components) of the electric and magnetic field intensities must be continuous at $z = 0$.

$$\vec{\mathbf{E}}_i(0) + \vec{\mathbf{E}}_r(0) = \vec{\mathbf{E}}_t(0) \quad (98)$$

$$\Rightarrow E_{i0} + E_{r0} = E_{t0} \quad (99)$$

$$\vec{\mathbf{H}}_i(0) + \vec{\mathbf{H}}_r(0) = \vec{\mathbf{H}}_t(0) \quad (100)$$

$$\Rightarrow \frac{E_{i0}}{\eta_1} - \frac{E_{r0}}{\eta_1} = \frac{E_{t0}}{\eta_2} \quad (101)$$

$$\begin{aligned} E_{i0} + E_{r0} &= E_{t0} \\ \frac{E_{i0}}{\eta_1} - \frac{E_{r0}}{\eta_1} &= \frac{E_{t0}}{\eta_2} \end{aligned} \quad (102)$$

$$\begin{aligned} E_{r0} - E_{t0} &= -E_{i0} \\ -\frac{E_{r0}}{\eta_1} - \frac{E_{t0}}{\eta_2} &= -\frac{E_{i0}}{\eta_1} \quad (// \eta_1) \end{aligned} \quad (103)$$

$$\begin{aligned} E_{r0} - E_{t0} &= -E_{i0} \\ -E_{r0} - \frac{\eta_1}{\eta_2} E_{t0} &= -E_{i0} \end{aligned} \quad (104)$$

$$-E_{t0} \left(1 + \frac{\eta_1}{\eta_2} \right) = -2 E_{i0} \quad (105)$$

$$E_{t0} \left(1 + \frac{\eta_1}{\eta_2} \right) = 2 E_{i0} \quad (106)$$

$$E_{t0} = \frac{2}{1 + \frac{\eta_1}{\eta_2}} E_{i0} = \frac{2}{\frac{\eta_2 + \eta_1}{\eta_2}} E_{i0} \quad (107)$$

$$E_{t0} = \frac{2 \eta_2}{\eta_1 + \eta_2} E_{i0} \quad (108)$$

$$E_{r0} = E_{t0} - E_{i0} \quad (109)$$

$$E_{r0} = \frac{2 \eta_2}{\eta_1 + \eta_2} E_{i0} - E_{i0} \quad (110)$$

$$E_{r0} = \left(\frac{2 \eta_2}{\eta_1 + \eta_2} - 1 \right) E_{i0} \quad (111)$$

$$E_{r0} = \frac{2 \eta_2 - (\eta_1 + \eta_2)}{\eta_1 + \eta_2} E_{i0} \quad (112)$$

$$E_{r0} = \frac{2 \eta_2 - \eta_1 - \eta_2}{\eta_1 + \eta_2} E_{i0} \quad (113)$$

$$E_{r0} = \frac{\eta_2 - \eta_1}{\eta_2 + \eta_1} E_{i0} \quad (114)$$

$$\Gamma = \frac{E_{r0}}{E_{i0}} = \frac{\eta_2 - \eta_1}{\eta_2 + \eta_1} \quad (\text{dimensionless}) \quad (\text{Reflection coefficient}) \quad (115)$$

$$\tau = \frac{E_{t0}}{E_{i0}} = \frac{2\eta_2}{\eta_2 + \eta_1} \quad (\text{dimensionless}) \quad (\text{Transmission coefficient}) \quad (116)$$

The reflection coefficient Γ can be positive or negative. The transmission coefficient τ is always positive.

$$E_{r0} = E_{t0} - E_{i0} \quad (117)$$

$$\frac{E_{r0}}{E_{i0}} = \frac{E_{t0}}{E_{i0}} - 1 \quad (118)$$

$$\Gamma = \tau - 1 \quad (119)$$

$$\tau = \Gamma + 1 \quad (120)$$

The total field in medium 1

$$\vec{E}_1(z) = \vec{E}_i(z) + \vec{E}_r(z) = \hat{a}_x E_{i0} (e^{-j\beta_1 z} + \Gamma e^{j\beta_1 z}) \quad (121)$$